## Equality condition in the

# data processing inequality for the quantum relative entropy 

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Let $P, Q$ be probability distributions on a discrete probability space $\mathcal{X}$, and define the Kullback-Leibler divergence $D_{\mathrm{KL}}(P \| Q)$ :

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D_{\mathrm{KL}}(P \| Q):=\sum_{x \in \mathcal{X}} P(x) \log \frac{P(x)}{Q(x)}
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This relative entropy is a premetric:

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Operational interpretation: Binary hypothesis testing

- Assume that we are given $n$ independent and identically distributed (i.i.d.) copies of one of two probability distributions $P$ or $Q$.
- Goal: Determine whether we have $P$ (null hypothesis $H_{P}$ )
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\begin{aligned}
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- In general: Trade-off between these errors.

One possibility: Try to minimize both at the same time
$\longrightarrow$ symmetric hypothesis testing, Chernoff bound

- Another one:
minimize type-II error s.t. type-I error $\leq \epsilon$
- Optimal exponent in the limit $n \rightarrow \infty$ given by $D_{\text {KL }}(P \| Q)$ :

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\text { type-ll error } \approx \exp \left(-n D_{K L}(P \| Q)\right) \text { forlarge } n \text {. }
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KL-divergence satisfies the data processing inequality (DPI):

- Let $P_{X}, Q_{X}$ be probability distributions on $\mathcal{X}$, and let $\Gamma_{Y \mid X}: X \rightarrow Y \in \mathcal{X}$ be a classical channel.
$\Rightarrow$ Denote by $P_{\gamma}, Q_{\gamma}$ the resulting distributions, that is, $P_{Y}(x):=\sum_{z \in \mathcal{X}} P_{X}(z) \Gamma_{Y \mid X}(z \mid x)$ and similar for $Q_{Y}$.
- Data processing inequality:

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D_{\text {KL }}\left(P_{X} \| Q_{X}\right) \geq D_{\text {KL }}\left(P_{Y} \| Q_{Y}\right)
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- Phrase an information-theoretic task in terms of transformations (e.g. encoding, decoding, ...).
- Characterize the task by entropic quantities based on relative entropies such as $D_{\text {KL }}(\cdot \| \cdot)$.
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How do we "make things quantum"?

- Replace the discrete probability space $\mathcal{X}$ by a Hilbert space $\mathcal{H}$ of dimension $|\mathcal{X}|<\infty$ (that is, $\mathcal{H} \cong \mathbb{C}^{|\mathcal{X}|}$ ). Density operator (or mixed state) is an operator $\rho$ acting on $\mathcal{H}$ that is
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- Interpretation: Assume that the pure state of a system is described by a normalized (column) vector $|\psi\rangle \in \mathcal{H}$.
$\Rightarrow$ Mixed state $\rho=\sum_{i} p_{i}\left|\psi_{i}\right\rangle\left\langle\psi_{i}\right|$ describes a system that is in the pure state $\psi_{i}$ with probability $p_{i}$. (in general, $\left|\mu_{i}\right\rangle$ \& $\left|\mu_{j}\right\rangle$ for $i \neq j$
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\rho=\sum_{i} \lambda_{i}\left|e_{i}\right\rangle\left\langle e_{i}\right| \quad \text { with }\left\langle e_{i} \mid e_{j}\right\rangle=\delta_{i j}
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where $\left|e_{i}\right\rangle$ is an eigenvector of $\rho$ with eigenvalue $\lambda_{i} \geq 0$.

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A bit more abstract:

- Density operators correspond to positive, normalized elements of the $C^{*}$-algebra $\mathcal{B}(\mathcal{H})$ of linear bounded operators acting on a Hilbert space $\mathcal{H}$ (for us $\operatorname{dim} \mathcal{H}<\infty$ ).
$\Rightarrow$ The $*-$ map is given by the adjoint ${ }^{\dagger}: A \mapsto A^{\dagger}$, and $\left\|A^{\dagger} A\right\|=\|A\|^{2}$ where $\|\cdot\|$ is the operator norm.
- Note that $\Lambda \geq 0 \Rightarrow \Lambda^{\dagger}=\Lambda$ (pos. elements are Hermitian). - We equip $\mathcal{B}(\mathcal{H})$ with the Hilbert-Schmidt inner product:

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\langle X, Y\rangle:=\operatorname{Tr}\left(X^{\dagger} Y\right)
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Dynamical evolution of a quantum system:

- A quantum channel (or quantum operation) is a map $\wedge: \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{K})$ that is

1 trace-preserving (TP): $\operatorname{Tr}(\Lambda(X))=\operatorname{Tr} X$ for all $X \in \mathcal{B}(\mathcal{H})$.
2 completely positive (CP): The map

is positive for all $n \in \mathbb{N}$.
$D$ Define the adjoint map $\wedge^{\dagger}: \mathcal{B}(K) \rightarrow \mathcal{B}(\mathcal{H})$ through

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Canonical example of quantum channel: Partial trace

- Consider two Hilbert spaces $\mathcal{H}_{1}, \mathcal{H}_{2}$.
- Define a linear map $\operatorname{Tr}_{1}: \mathcal{B}\left(\mathcal{H}_{1} \otimes \mathcal{H}_{2}\right) \rightarrow \mathcal{B}\left(\mathcal{H}_{2}\right)$ by

$$
\operatorname{Tr}_{1}(x \otimes Y)=\operatorname{Tr}(x) Y
$$

for arbitrary $X \in \mathcal{B}\left(\mathcal{H}_{1}\right), Y \in \mathcal{B}\left(\mathcal{H}_{2}\right)$.

- Trace-preserving: $\operatorname{Tr}\left(\operatorname{Tr}_{1}(X \otimes Y)\right)=\operatorname{Tr}(X) \operatorname{Tr}(Y)=\operatorname{Tr}(X \otimes Y)$
- Completely positive: $\operatorname{Tr}_{1} \otimes \mathrm{id}_{n}$ is positive for all $n \in \mathbb{N}$.


## Mathematics of Quantum Mechanics 101

Canonical example of quantum channel: Partial trace

- Consider two Hilbert spaces $\mathcal{H}_{1}, \mathcal{H}_{2}$.
- Define a linear map $\operatorname{Tr}_{1}: \mathcal{B}\left(\mathcal{H}_{1} \otimes \mathcal{H}_{2}\right) \rightarrow \mathcal{B}\left(\mathcal{H}_{2}\right)$ by

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\operatorname{Tr}_{1}(X \otimes Y)=\operatorname{Tr}(X) Y
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for arbitrary $X \in \mathcal{B}\left(\mathcal{H}_{1}\right), Y \in \mathcal{B}\left(\mathcal{H}_{2}\right)$.

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Importance of partial trace

1 Stinespring's representation theorem [Stinespring 1955]:
Let $\Lambda: \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{K})$ be a quantum channel, then there is an isometry $V: \mathcal{H} \rightarrow \mathcal{K} \otimes \mathcal{K}^{\prime}$ s.t. $\Lambda(X)=\operatorname{Tr}_{2}\left(V X V^{\dagger}\right)$.
(Every quantum ch. looks like the partial trace in some space.)

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## Mathematics of Quantum Mechanics 101

Operators and functional calculus

- We write $A \geq B$ if $A-B \geq 0$.
$\Rightarrow$ Let $A \geq 0$ with spectral decomposition $A=\sum_{i} \lambda_{i}\left|e_{i}\right\rangle\left\langle e_{i}\right|$,
and let $f: \mathbb{R}^{+} \rightarrow \mathbb{R}$, then we define $f(A):=\sum_{i} f\left(\lambda_{i}\right)\left|e_{i}\right\rangle\left\langle e_{i}\right|$
- $f$ is operator monotone: $A \geq B$ implies $f(A) \geq f(B)$.
$\rightarrow f$ is operator convex: For $\lambda \in(0,1)$ and $A_{1}, A_{2} \geq 0$,

$$
f\left(\lambda A_{1}+(1-\lambda) A_{2}\right) \leq \lambda f\left(A_{1}\right)+(1-\lambda) f\left(A_{2}\right)
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- Jensen's operator inequality: $f$ is operator convex iff

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## Quantum relative entropy

Let $X, Y \in \mathcal{B}(\mathcal{H}), X, Y \geq 0$ with $\operatorname{supp} X \subseteq \operatorname{supp} Y$, then

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D(X \| Y):=\operatorname{Tr}[X(\log X-\log Y)]
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- If $X$ and $Y$ are states, then $D(X \| Y) \geq 0$, and $=0$ iff $X=Y$.
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$\triangleright D(\hat{P} \| \hat{Q})=D_{\text {KL }}(P \| Q)$ for classical states

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error exponent in quantum hypothesis testing
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- Let $\Phi: \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{K})$ be TP and 2-positive map.

2-positive: For $A_{i} \in \mathcal{B}(\mathcal{H}), i=1, \ldots, 4$ we have

$$
\left(\begin{array}{ll}
\Phi\left(A_{1}\right) & \Phi\left(A_{2}\right) \\
\Phi\left(A_{3}\right) & \Phi\left(A_{4}\right)
\end{array}\right) \geq 0 \quad \text { if } \quad\left(\begin{array}{ll}
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- Then $D(\cdot \| \cdot)$ satisfies the data processing inequality $D(X \| Y) \geq D(\Phi(X) \| \Phi(Y))$.
- This holds in particular for every CPTP map $\wedge$. Recall: $\Phi$ CP : $\Leftrightarrow \Phi$ is $n$-positive for all $n$
- Recent result: DPI holds for every positive map.
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## Equality in data processing inequality

Theorem (Petz 1988)
Let $\Phi: \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{K})$ be a 2-positive TP map, and let
$X, Y \in \mathcal{B}(\mathcal{H})$ be invertible density operators. Then we have

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if and only if for all $t \in \mathbb{R}$

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- Algebraic condition on map and operators.
- Equivalent formulation: There exists a recovery map

such that $\mathcal{R}_{\Phi, Y}(\Phi(X))=X$ and $\mathcal{R}_{\Phi, Y}(\Phi(Y))=Y$.
- $\mathcal{R}_{\Phi}$, recovers $X, Y$ by reverting the action of $\Phi$.
- Recovery on $X$ :

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- Consider the multiplication operators $L_{A}(T):=A T$ and $R_{D}(T)=T B$, satisfying


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$L_{A^{-1}}=L_{A}^{-1}$ if $A$ is invertible, likewise for $R_{B}$.
$L_{A}, R_{B}$ are self-adjoint, and positive if $A, B \geq 0$.
For analytic $f: \mathbb{R}^{+} \rightarrow \mathbb{R}$ we have $f\left(L_{A}\right)=L_{f(A)}$ and likewise
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- Define the relative modular operator $\Delta_{Y \mid X}=L_{Y} \circ R_{X^{-1}}$.

$$
\begin{aligned}
& \text { Then } \log \Delta_{Y \mid X}=L_{\log Y}-R_{\log X} \text {, and } \\
& \qquad D(X \| Y)=\operatorname{Tr}[X(\log X-\log Y)]=-\left\langle X^{1 / 2}, \log \Delta_{Y \mid X}\left(X^{1 / 2}\right)\right\rangle . \\
& \text { Assume now that } \Phi(X) \text { is also invertible, and set } \\
& \qquad \Delta \equiv \Delta_{Y \mid X} \quad \Delta_{\Phi} \equiv \Delta_{\Phi(Y) \mid \Phi(X)}
\end{aligned}
$$

such that

$$
D(X \| Y)=-\left\langle X^{1 / 2}, \log \Delta\left(X^{1 / 2}\right)\right\rangle
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$$
\Delta \equiv \Delta_{Y \mid X} \quad \Delta_{\Phi} \equiv \Delta_{\Phi(Y) \mid \Phi(X)}
$$

such that

$$
\begin{aligned}
D(X \| Y) & =-\left\langle X^{1 / 2}, \log \Delta\left(X^{1 / 2}\right)\right\rangle \\
D(\Phi(X) \| \Phi(Y)) & =-\left\langle\Phi(X)^{1 / 2}, \log \Delta_{\Phi}\left(\Phi(X)^{1 / 2}\right)\right\rangle
\end{aligned}
$$

## Proof of the main theorem

- Consider the integral representation

$$
\log x=\int_{0}^{\infty} \frac{1}{1+t}-\frac{1}{x+t} d t
$$

- We can then write


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-(1+t)^{-1} d t
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and focus on the integrands written in terms of the
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- Define a linear map $V: \mathcal{B}(\mathcal{K}) \rightarrow \mathcal{B}(\mathcal{H})$ :

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V(A):=\Phi^{\dagger}\left(A \Phi(X)^{-1 / 2}\right) X^{1 / 2} .
$$

$\Phi^{\dagger}$ is unital: $V\left(\Phi(X)^{1 / 2}\right)=\Phi^{\dagger}(\mathbb{1}) X^{1 / 2}=X^{1 / 2}$.

- $V$ is a contraction: $\|V(A)\|^{2} \leq\|A\|$ for all $A$.
- To show this, use the Schwarz inequality

$$
\Phi^{\dagger}\left(A^{\dagger} A\right) \geq \Phi^{\dagger}\left(A^{\dagger}\right) \Phi^{\dagger}(A) .
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2-positive TP maps largest class of maps for which SI holds!

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again by application of the Schwarz inequality.

## Proof of the main theorem

- $V$ is a contraction, and $V^{\dagger} \Delta V \leq \Delta_{\Phi}$.
$\rightarrow y \mapsto(y+t)^{-1}$ is operator monotone (decreasing) and
operator convex:

$$
\left(\Delta_{\Phi}+t\right)^{-1} \leq\left(V^{\dagger} \Delta V+t\right)^{-1} \leq V^{\dagger}(\Delta+t)^{-1} V
$$

(second $\leq$ follows from Jensen's operator inequality)

- Hence:
$\left\langle X^{1 / 2},(\Delta+t)^{-1}\left(X^{1 / 2}\right)\right\rangle=\left\langle V \Phi(X)^{1 / 2},(\Delta+t)^{1 / 2} V\left(\Phi(X)^{1 / 2}\right)\right\rangle$

$$
\begin{aligned}
& =\left\langle\Phi(x)^{1 / 2} V^{\dagger}(\Delta+t)^{1 / 2} V\left(\Phi(x)^{1 / 2}\right)\right\rangle \\
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- Recall integral representations:

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- Just proved:

- Insert this in the integral representations:



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- Equality in DPI if and only if for all $t>0$

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Let $A \geq B$, then $\langle v| A|v\rangle=\langle v| B|v\rangle$ implies $A|v\rangle=B|v\rangle$. - Hence, $V^{\dagger}(\Delta+t)^{-1} V\left(\Phi(X)^{1 / 2}\right)=\left(\Delta_{\Phi}+t\right)^{-1}\left(\Phi(X)^{1 / 2}\right)$. - It follows by an easy calculation that


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\left\|V^{\dagger}(\Delta+t)^{-1}\left(X^{1 / 2}\right)\right\|^{2} & =\left\|V^{\dagger}(\Delta+t)^{-1} V\left(\Phi(X)^{1 / 2}\right)\right\|^{2} \\
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- $\left\|V^{\dagger}(\Delta+t)^{-1}\left(X^{1 / 2}\right)\right\|=\left\|(\Delta+t)^{-1}\left(X^{1 / 2}\right)\right\|^{2}$.
- If $\left.\| W^{\dagger}|\xi\rangle\left\|^{2}=\right\| \| \xi\right\rangle \|^{2}$ for an arbitrary contraction $W$, then $W W^{\dagger}|\xi\rangle=|\xi\rangle$
- Recall: $V^{\dagger}(\Delta+t)^{-1}\left(X^{1 / 2}\right)=\left(\Delta_{\Phi}+t\right)^{-1}\left(\Phi(X)^{1 / 2}\right)$
- Then: $V\left(\Delta_{\Phi}+t\right)^{-1}\left(\Phi(X)^{1 / 2}\right)=(\Delta+t)^{-1}\left(X^{1 / 2}\right)$.
- That is, the resolvents of $\Delta_{\Phi}$ and $\Delta$ coincide on the vectors $\Phi(X)^{1 / 2}$ and $X^{1 / 2}$, respectivelv (modulo $V$ ).


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## Proof of the main theorem

- The resolvent of an operator $O$ determines the projections onto the eigenspaces of $O$.
- Hence, for every polynomial $p$ we have

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V p\left(\Delta_{\Phi}\right)\left(\Phi(X)^{1 / 2}\right)=p(\Delta)\left(X^{1 / 2}\right)
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- Stone-Weierstrass approximation theorem: polynomials are dense in the space of continuous functions
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## Proof of the main theorem

- $V f\left(\Delta_{\Phi}\right)\left(\Phi(X)^{1 / 2}\right)=f(\Delta)\left(X^{1 / 2}\right)$ for all continuous $f$. - Choose $f(x)=x^{\text {it }}$, then



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The last line follows from taking the adjoint and using the fact that $\wedge^{\prime}\left(A^{+}=\wedge^{( } \wedge^{+}\right)$for a positive map $\wedge$.

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\Leftrightarrow V\left(\Phi(Y)^{i t} \Phi(X)^{-i t} \Phi(X)^{1 / 2}\right) & =Y^{i t} X^{-i t} X^{1 / 2} \\
\Leftrightarrow \Phi^{\dagger}\left(\Phi(Y)^{i t} \Phi(X)^{-i t}\right) X^{1 / 2} & =Y^{i t} X^{-i t} X^{1 / 2} \\
\Leftrightarrow \Phi^{\dagger}\left(\Phi(Y)^{i t} \Phi(X)^{-i t}\right) & =Y^{i t} X^{-i t} \\
\Leftrightarrow \Phi^{\dagger}\left(\Phi(X)^{i t} \Phi(Y)^{-i t}\right) & =X^{i t} Y^{-i t}
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\end{aligned}
$$

The last line follows from taking the adjoint and using the fact that $\Lambda(A)^{\dagger}=\Lambda\left(A^{\dagger}\right)$ for a positive map $\Lambda$.

## Proof of the main theorem

- Proven so far:

$$
\begin{aligned}
D(X \| Y) & =D(\Phi(X) \| \Phi(Y)) \\
& \Longrightarrow \Phi^{\dagger}\left(\Phi(X)^{i t} \Phi(Y)^{-i t}\right)=X^{i t} Y^{-i t} \quad \text { for all } t>0
\end{aligned}
$$

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## von Neumann entropy

For a state $\rho$ define the von Neumann entropy

$$
S(\rho)=-\operatorname{Tr} \rho \log \rho=-D(\rho \| \mathbb{1})
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- $S(\rho)=H\left(\left\{\lambda_{i}\right\}_{i}\right)$ where $\lambda_{i}$ are the eigenvalues of $\rho$ and $H\left(\left\{p_{i}\right\}_{i}\right)=-\sum_{i} p_{i} \log p_{i}$ is the Shannon entropy.
$>0 \leq S(\rho) \leq \log \operatorname{dim} \mathcal{H}$ for all states $\rho$ on $\mathcal{H}$.
- Additivity: $S(\rho \otimes \sigma)=S(\rho)+S(\sigma)$
- Subadditivity: Let $\rho_{A B}$ be a state on a bipartite system $\mathcal{H}_{A} \otimes \mathcal{H}_{B}$ and set $\rho_{A}=\operatorname{Tr}_{B} \rho_{A B}$ and $\rho_{B}=\operatorname{Tr}_{A} \rho_{A B}$. Then,

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$\Rightarrow$ Additivity: $S(\rho \otimes \sigma)=S(\rho)+S(\sigma)$
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## Strong subadditivity of von Neumann entropy

Strong subadditivity [Lieb and Ruskai 1973]:

- Let $\rho_{A B C}$ be a tripartite state, and denote by $\rho_{A B}, \rho_{B C}, \rho_{B}$ the corresponding marginals. Then:

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I(A ; C \mid B)_{\rho}=S(A B)+S(B C)-S(B)-S(A B C)
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where $S(A B) \equiv S\left(\rho_{A B}\right)$ etc.

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## Strong subadditivity of von Neumann entropy

Equality in SSA [Hayden et al. 2004]:

- We applied DPI with respect to partial trace over $C$.
- Equality condition (recovery map formulation):

There is a recovery map $\mathcal{R}_{B \rightarrow B C}: \mathcal{B}\left(\mathcal{H}_{B}\right) \rightarrow \mathcal{B}\left(\mathcal{H}_{B C}\right)$ s.t.

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- Hence, we obtain:
$I(A ; C \mid B)_{\rho}=0 \longleftrightarrow \exists \mathcal{R}_{B \rightarrow B C}$ with $\rho_{B C}=\mathcal{R}_{B \rightarrow B C}\left(\rho_{B}\right)$
- That is, $A \leftrightarrow B \leftrightarrow C$ forms a (short) quantum Markov
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## Summary:

- Divergences (or relative entropies) play an important role in Classical and Quantum Information Theory.
- Their crucial property is the data processing inequality. The quantum relative entropy is an important example in Quantum Information Theory
- We derived an equality condition in the DPI for the quantum relative entropy.
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Generalized divergence:

- Crucial property of a divergence: DPI

$>\lim _{\alpha \rightarrow 1} D_{\alpha}(\rho \| \sigma)=D(\rho \| \sigma)=\lim _{\alpha \rightarrow 1} \widetilde{D}_{\alpha}(\rho \| \sigma)$.
- Both satisfy DPI in the given range.
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Thank you very much for your attention!

